# Aspects of string phenomenology in particle physics and cosmology

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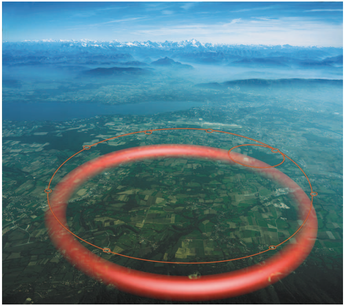
# String phenomenology

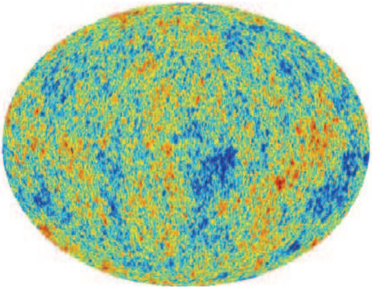
- Is string theory a tool for strong coupling dynamics

  or a theory of fundamental forces?

- If theory of Nature can string theory describe

  both particle physics and cosmology?





# Problem of scales

- describe high energy SUSY extension of the Standard Model

  unification of all fundamental interactions

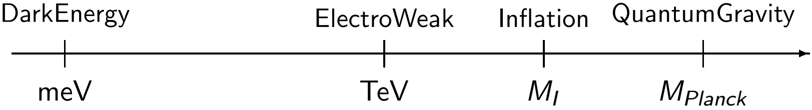
- incorporate Dark Energy

  simplest case: infinitesimal (tunable) +ve cosmological constant

- describe possible accelerated expanding phase of our universe

  models of inflation (approximate de Sitter)

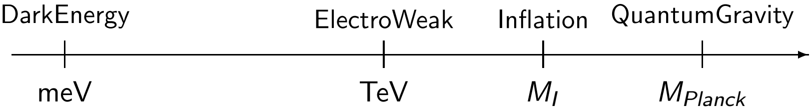
⇒ 3 very different scales besides $M_{Planck}$ :



DarkEnergy ElectroWeak Inflation QuantumGravity

meV TeV $M_I$ $M_{Planck}$

# Problem of scales


DarkEnergy ElectroWeak Inflation QuantumGravity

meV TeV $M_I$ $M_{Planck}$

1. possible connections
   - $M_I$ could be near the EW scale, such as in Higgs inflation

     but large non minimal coupling to explain
   - $M_{Planck}$ could be emergent from the EW scale

     in models of low-scale gravity and TeV strings

     2 extra dims at submm $\leftrightarrow$ meV: interesting coincidence with DE scale

     $M_I \sim TeV$ is also allowed by the data since cosmological observables are dimensionless in units of the effective gravity scale

     I.A.-Patil '14
2. they are independent [8]

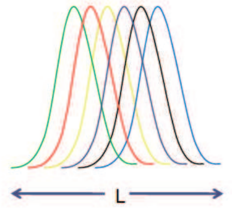
# Effective scale of gravity: reduced by the number of species

$N$ particle species $\Rightarrow$ lower quantum gravity scale : $M_*^2 = M_p^2/N$

Dvali '07, Dvali, Redi, Brustein, Veneziano, Gomez, Lüst '07-'10

derivation from: black hole evaporation or quantum information storage

Pixel of size $L$ containing $N$ species storing information:



localization energy $E \gtrsim N/L \rightarrow$

Schwarzschild radius $R_s = N/(LM_p^2)$

no collapse to a black hole : $L \gtrsim R_s \Rightarrow L \gtrsim \sqrt{N}/M_p = 1/M_*$

# Cosmological observables

Power spectrum of temperature anisotropies

(adiabatic curvature perturbations $\mathcal{R}$)

$$\mathcal{P}_{\mathcal{R}} = \frac{H^2}{8\pi^2 M_*^2 \epsilon} \simeq \mathcal{A} \times 10^{-10} \quad ; \quad \mathcal{A} \approx 22$$

$$\nwarrow -\dot{H}/H^2$$

Power spectrum of primordial tensor anisotropies $\mathcal{P}_t = 2\frac{H^2}{\pi^2 M_*^2}$

$\Rightarrow$ tensor to scalar ratio $r = \mathcal{P}_t/\mathcal{P}_{\mathcal{R}} = 16\epsilon$

measurement of $\mathcal{A}$ and $r$ $\Rightarrow$ fix the scale of inflation

$$H \text{ in terms of } M_* \quad : \quad \frac{H}{M_*} = \left(\frac{\pi^2 \mathcal{A}\, r}{2 \times 10^{10}}\right)^{1/2} \equiv \Upsilon \approx 1.05\sqrt{r} \times 10^{-4}$$

# Extra species as Kaluza-Klein states

$D = 4 + n$ extra dims of size average size $R \Rrightarrow$

fundamental gravity scale $M_{**}^{2+n} R^n = M_{Pl}^2$

$N =$ all KK states with mass less than $H \Rrightarrow N \simeq (HR)^n$

$M_* = M_{Pl}/\sqrt{N} = M_{**}(M_{**}R)^{n/2}/(HR)^{n/2} = M_{**}(M_{**}/H)^{n/2}$

$H = M_* \Upsilon = M_{**}(M_{**}/H)^{n/2}\Upsilon \quad \Rrightarrow \quad H = M_{**}\Upsilon^{2/(n+2)}$

$\Rrightarrow H \sim$ 1-3 orders of magnitude less than $M_{**}$ for $0.001 \lesssim r \lesssim 0.1$

as low as near the EW scale [4]

# impose independent scales: proceed in 2 steps

1. SUSY breaking at $m_{SUSY} \sim$ TeV

   with an infinitesimal (tunable) positive cosmological constant

   Villadoro-Zwirner '05

   I.A.-Knoops, I.A.-Ghilencea-Knoops '14, I.A.-Knoops in preparation

2. Inflation in supergravity at a scale different than $m_{SUSY}$ [27]

1st step: Maximal predictive power if there is common framework for :

- moduli stabilization
- model building (spectrum and couplings)
- SUSY breaking (calculable soft terms)
- computable radiative corrections (crucial for comparing models)

Possible candidate of such a framework: **magnetized branes**

# Type I string theory with magnetic fluxes $B_{ij}$ on 2-cycles of the compactification manifold

- Dirac quantization: $B = \frac{m}{nA} \equiv \frac{p}{A}$ [12] $\Rightarrow$ moduli stabilization

  $B$: constant magnetic field    $m$: units of magnetic flux

  $n$: brane wrapping    $A$: area of the 2-cycle

- Spin-dependent mass shifts for charged states $\Rightarrow$ SUSY breaking
- Exact open string description: $\Rightarrow$ calculability

  $qB \to \theta = \arctan qB\alpha'$    weak field $\Rightarrow$ field theory

- T-dual representation: branes at angles $\Rightarrow$ model building

  $(m, n)$: wrapping numbers around the 2-cycle directions

# Magnetic fluxes can be used to stabilize moduli

**I.A.-Maillard '04, I.A.-Kumar-Maillard '05, '06, Bianchi-Trevigne '05**

e.g. $T^6$: 36 moduli (geometric deformations)

internal metric: $6 \times 7/2 = 21 = 9 + 2 \times 6$

type IIB RR 2-form: $6 \times 5/2 = 15 = 9 + 2 \times 3$

$$\text{complexification} \Rightarrow \begin{cases} \text{Kähler class} & J \\ \\ \text{complex structure} & \tau \end{cases} \qquad \text{9 complex moduli for each}$$

magnetic flux: $6 \times 6$ antisymmetric matrix $F$ complexification $\Rightarrow$

$F_{(2,0)}$ on holomorphic 2-cycles: potential for $\tau$ superpotential

$F_{(1,1)}$ on mixed (1,1)-cycles: potential for $J$ FI D-terms

# $N = 1$ SUSY conditions ⇒ moduli stabilization

1. $F_{(2,0)} = 0 \;\Rightarrow \tau$ matrix equation for every magnetized $U(1)$

   need 'oblique' (non-commuting) magnetic fields to fix off-diagonal components of the metric ← but can be made diagonal

2. $J \wedge J \wedge F_{(1,1)} = F_{(1,1)} \wedge F_{(1,1)} \wedge F_{(1,1)} \;\Rightarrow J$

   vanishing of a Fayet-Iliopoulos term: $\xi \sim F \wedge F \wedge F - J \wedge J \wedge F$

   magnetized $U(1)$ → massive absorbs RR axion

   one condition ⇒ need at least 9 brane stacks

3. Tadpole cancellation conditions : introduce an extra brane(s)

   ⇒ dilaton potential from the FI D-term → two possibilities:

   - keep SUSY by turning on charged scalar VEVs
   - break SUSY in a dS or AdS vacuum $d = \xi/\sqrt{1+\xi^2}$ [13]

I.A.-Derendinger-Maillard '08

$F_{(2,0)} = 0 \Rightarrow \qquad \tau^{\mathrm{T}} p_{xx} \tau - (\tau^{\mathrm{T}} p_{xy} + p_{yx} \tau) + p_{yy} = 0$ [9]

$T^6$ parametrization: $(x^i, y^i) \quad i = 1, 2, 3 \qquad z^i = x^i + \tau^{ij} y^i$

Non-trivial VEVs $v$ for charged brane scalars $\Rightarrow$

D-term condition is modified to:

$$q\, v^2 \,(J \wedge J \wedge J - J \wedge F \wedge F) = -(F \wedge F \wedge F - F \wedge J \wedge J)$$

charge

# break SUSY in a dS or AdS vacuum

**I.A.-Derendinger-Maillard ’08**

$N = 2$ non-linear supersymmetry $\Rightarrow$

General form of the localized dilaton potential:

$$V(\phi, d) = \frac{e^{-\phi}}{g^2}\left\{\left(\sqrt{1-d^2}-1\right) + \xi d + \delta T\right\}$$

DBI action FI-term

$d$: D-auxiliary in $2\pi\alpha'$-units

$\delta T$: tension leftover RR tadpole cancellation $\Rightarrow \delta T = 1 - \sqrt{1-\xi^2}$

$d$ elimination $\Rightarrow$ $d = \frac{\xi}{\sqrt{1+\xi^2}}$

$$V_{\min} = \delta\bar{T} e^{-\phi} \quad ; \quad \delta\bar{T} = \sqrt{1+\xi^2} - \sqrt{1-\xi^2}$$

# Dilaton fixing:

1) by 3-form fluxes in a SUSY way $\Rightarrow$ dS vacuum with positive energy

D-term uplifting possible from flat space

2) add a 'non-critical' (bulk) dilaton potential

$\Rightarrow$ AdS vacuum with tunable string coupling

$V_{\mathrm{non-crit}} = \delta c \, e^{-2\phi}$ $\delta c$: central charge deficit

minimization of $V = V_{\mathrm{non-crit}} + V_{\mathrm{min}} \Rightarrow \delta c < 0$

$e^{\phi_0} = -\frac{2\delta c}{3\delta \bar{T}}$ $V_0 = \frac{\delta c^3}{3\delta \bar{T}^2}$ $R_0 = -\delta \bar{T} e^{3\phi_0}$

curvature in Einstein frame

e.g. replace a free coordinate by a CFT minimal model of central charge $1 + \delta c$

→ generalize: add a dilaton potential preserving the axion shift symmetry

$\Rightarrow$ break SUSY with tunable vacuum energy

# Toy model for SUSY breaking

Content (besides $N = 1$ SUGRA): one vector $V$ and one chiral multiplet $S$

with a shift symmetry $S \to S - ic\omega \leftarrow$ transfomration parameter

String theory: compactification modulus or universal dilaton

$s = 1/g^2 + ia \leftarrow$ dual to antisymmetric tensor

Kähler potential $K$: function of $S + \bar{S}$

string theory: $K = -p \ln(S + \bar{S})$

Superpotential: constant or single exponential if R-symmetry $W = ae^{bS}$

$b < 0 \Rightarrow$ non perturbative

can also be described by a generalized linear multiplet

# Scalar potential

$$\mathcal{V}_F = a^2 e^{\frac{b}{l}} l^{p-2} \left\{ \frac{1}{p}(pl-b)^2 - 3l^2 \right\} \qquad l = 1/(s+\bar{s})$$

Planck units

no minimum for $b < 0$ with $l > 0$ $(p \leq 3)$

but interesting metastable SUSY breaking vacuum

when R-symmetry is gauged by $V$ allowing a Fayet-Iliopoulos (FI) term:

$$\mathcal{V}_D = c^2 l(pl-b)^2$$

for gauge kinetic function $f(S) = S$

- $b > 0$: $\mathcal{V} = \mathcal{V}_F + \mathcal{V}_D$ SUSY local minimum in AdS space at $l = b/p$
- $b = 0$: SUSY breaking minimum in AdS $(p < 3)$ $\delta c = -a^2$
- $b < 0$: SUSY breaking minimum with tunable cosmological constant $\Lambda$

In the limit $\Lambda \approx 0$ $(p = 2)$ $\Rightarrow$

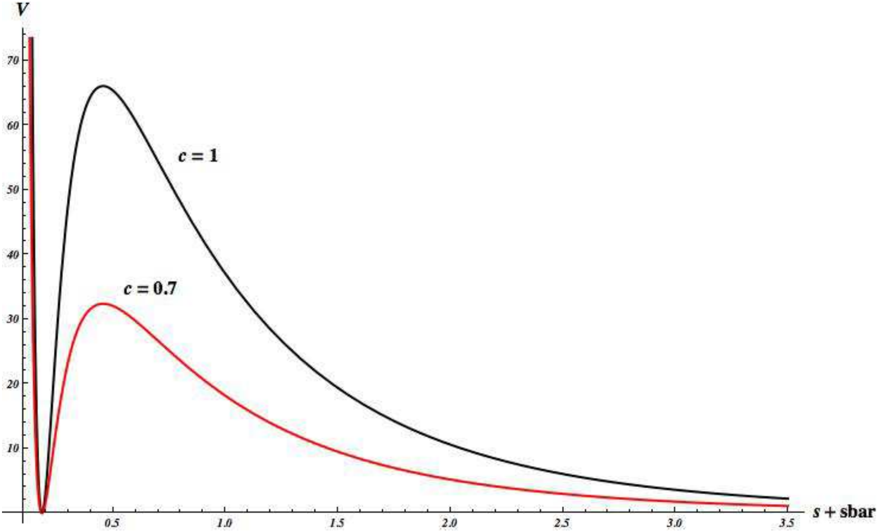
$b/l = \alpha \approx -0.183268$

$\frac{a^2}{bc^2} = 2\frac{e^{-\alpha}}{\alpha}\frac{(2-\alpha)^2}{2+4\alpha-\alpha^2} + \mathcal{O}(\Lambda) \approx -50.6602$

physical spectrum:

massive dilaton, $U(1)$ gauge field, Majorana fermion, gravitino

All masses of order $m_{3/2} \approx e^{\alpha/2} l a \leftarrow$ TeV scale


V
70
60
50
40
30
20
10
$c = 1$
$c = 0.7$
0.5
1.0
1.5
2.0
2.5
3.0
3.5
s + sbar

# Properties and generalizations

- Metastability of the ground state: extremely long lived

  $l \simeq 0.02$ (GUT value $\alpha_{GUT}/2$) $m_{3/2} \sim \mathcal{O}(TeV) \Rrightarrow$

  decay rate $\Gamma \sim e^{-\mathcal{B}}$ with $B \approx 10^{300}$

- Add visible sector (MSSM) preserving the same vacuum

  matter fields $\phi$ neutral under R-symmetry

  $$K = -2\ln(S+\bar{S}) + \phi^\dagger\phi \quad ; \quad W = (a + W_{MSSM})e^{bS}$$

  $\Rrightarrow$ soft scalar masses non-tachyonic of order $m_{3/2}$ (gravity mediation)

- R-charged fields can be added in the hidden sector

  needed for anomaly cancellation (important constraint)

# Properties and generalizations

- Interesting phenomenology: work in progress

- Toy model classically equivalent to

  $K = -p\ln(S+\bar{S}) + b(S+\bar{S}) \quad ; \quad W = a$ with $V$ ordinary $U(1)$

- string origin of $b$ ? allows flat space solution

  unphysical in the absence of $a$

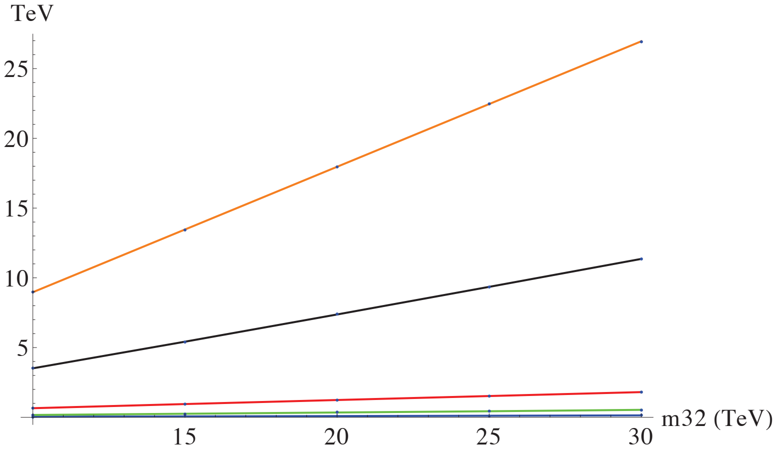
- Consider a simple (anomaly free) variation of the model with

  the above $K$ and $W$, gauge kinetic function $f = 1$ and $p = 1$

  $\Rightarrow$ tuning still possible but scalar masses of neutral matter tachyonic

  possible solution: add a new field $Z$ in the 'hidden' $\mathrm{S}\not{\mathrm{U}}\mathrm{SY}$ sector [8]

# Typical spectrum



The masses of sbottom squark (yellow), stop (black), gluino (red), lightest chargino (green) and lightest neutralino (blue) as a function of the gravitino mass. The mass of the lightest neutralino varies between 42 GeV for $m_{3/2}$=10 TeV and 138 GeV for $m_{3/2}$=30 TeV.

# Starobinsky model of inflation

$$\mathcal{L} = \frac{1}{2}R + \alpha R^2$$

Lagrange multiplier $\phi \Rightarrow \mathcal{L} = \frac{1}{2}(1+2\phi)R - \frac{1}{4\alpha}\phi^2$

Weyl rescaling $\Rightarrow$ equivalent to a scalar field with exponential potential:

$$\mathcal{L} = \frac{1}{2}R - \frac{1}{2}(\partial\phi)^2 - \frac{M^2}{12}\left(1 - e^{-\sqrt{\frac{2}{3}}\phi}\right)^2 \qquad M^2 = \frac{3}{4\alpha}$$

Note that the two metrics are not the same

supersymmetric extension:

add D-term $\mathcal{R}\bar{\mathcal{R}}$ because F-term $\mathcal{R}^2$ does not contain $R^2$

$\Rightarrow$ brings two chiral multiplets

# SUSY extension of Starobinsky model

$$K = -3\ln(T + \bar{T} - C\bar{C}) \quad ; \quad W = MC(T - \frac{1}{2})$$

- $T$ contains the inflaton: $\mathrm{Re}\,T = e^{\sqrt{\frac{2}{3}}\phi}$
- $C \sim \mathcal{R}$ is unstable during inflation

  $\Rightarrow$ add higher order terms to stabilize it

  e.g. $C\bar{C} \to h(C, \bar{C}) = C\bar{C} - \zeta(C\bar{C})^2$ Kallosh-Linde '13
- SUSY is broken during inflation with $C$ the goldstino superfield [33]

# Non-linear supersymmetry $\Rightarrow$ goldstino mode $\chi$

**Volkov-Akulov '73**

Why study goldstino interactions:

- Effective field theory of SUSY breaking at low energies $m_\chi << m_{susy}$

  e.g. gauge mediation dominant vs gravity mediation

  $\chi$: longitudinal gravitino with $m_\chi \simeq \frac{m^2_{susy}}{M_{Planck}} \lesssim m_{soft} << m_{susy}$

  $M_{Planck} \to \infty$: SUGRA decoupled

  massless $\chi$ coupled to matter $\sim 1/m_{susy}$

- Brane dynamics: half SUSY of the bulk broken but NL realized

Non-linear SUSY transformations:

$$\delta\chi_\alpha = \frac{\xi_\alpha}{\kappa} + \kappa\,\Lambda^\mu_\xi\,\partial_\mu\chi_\alpha \qquad \Lambda^\mu_\xi = -i\left(\chi\sigma^\mu\bar{\xi} - \xi\sigma^\mu\bar{\chi}\right)$$

$\kappa$: goldstino decay constant (SUSY breaking scale) $\qquad \kappa = (\sqrt{2}m_{susy})^{-2}$

**Volkov-Akulov action:**

Define the ‘vierbein’: $\quad E^a_\mu = \delta^a_\mu + \kappa^2\, t^a_\mu \qquad t^a_\mu = i\chi\overset{\leftrightarrow}{\partial}_\mu\sigma^a\bar{\chi}$

$\delta(\det E) = \kappa\,\partial_\mu\left(\Lambda^\mu_\xi\det E\right) \Rightarrow$ invariant action:

$$S_{VA} = -\tfrac{1}{2\kappa^2}\int d^4x\,\det E = -\tfrac{1}{2\kappa^2} - \tfrac{i}{2}\chi\sigma^\mu\overset{\leftrightarrow}{\partial}_\mu\bar{\chi} + \dots$$

# Constrained superfields

**Rocek-Tseytlin '78, Lindstrom-Rocek '79, Komargodski-Seiberg '09**

spontaneous global $\text{S}\!\!\!/\text{USY}$ : no supercharge but still conserved supercurrent

$\Rightarrow$ superpartners exist in operator space (not as 1-particle states)

$\Rightarrow$ constrained superfields: 'eliminate' superpartners

Goldstino: chiral superfield $X_{NL}$ satisfying $X_{NL}^2 = 0 \Rightarrow$

$$X_{NL}(y) = \frac{\chi^2}{2F} + \sqrt{2}\theta\chi + \theta^2 F \qquad y^\mu = x^\mu + i\theta\sigma^\mu\bar{\theta}$$

$$= F\Theta^2 \qquad \Theta = \theta + \frac{\chi}{\sqrt{2}F}$$

$$\mathcal{L}_{NL} = \int d^4\theta X_{NL}\bar{X}_{NL} - \frac{1}{\sqrt{2}\kappa}\left\{\int d^2\theta X_{NL} + h.c.\right\} = \mathcal{L}_{Volkov-Akulov}$$

$$F = \frac{1}{\sqrt{2}\kappa} + \ldots$$

# Non-linear SUSY in supergravity

**I.A.-Dudas-Ferrara-Sagnotti ’14**

$$K = -3\log(1 - X\bar{X}) \equiv 3X\bar{X} \quad ; \quad W = f\,X + W_0 \qquad X \equiv X_{NL}$$

$$\Rightarrow \quad V = \frac{1}{3}|f|^2 - 3|W_0|^2 \quad ; \quad m_{3/2}^2 = |W_0|^2$$

- $V$ can have any sign   contrary to global NL SUSY
- NL SUSY in flat space $\Rightarrow f = 3\,m_{3/2}M_p$
- Dual gravitational formulation: $(\mathcal{R} - 6W_0)^2 = 0$   I.A.-Markou ’15

  chiral curvature superfield
- Minimal SUSY extension of $R^2$ gravity [28]

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Minimal SUSY extension that evades stability problem:

replace $C$ by the non-linear multiplet $X$

# Non-linear Starobinsky supergravity

$$K = -3\ln(T + \bar{T} - X\bar{X}) \quad ; \quad W = M\,XT + f\,X + W_0 \qquad \Rightarrow$$

$$\mathcal{L} = \frac{1}{2}R - \frac{1}{2}(\partial\phi)^2 - \frac{M^2}{12}\left(1 - e^{-\sqrt{\frac{2}{3}}\phi}\right)^2 - \frac{1}{2}e^{-2\sqrt{\frac{2}{3}}\phi}(\partial a)^2 - \frac{M^2}{18}e^{-2\sqrt{\frac{2}{3}}\phi}a^2$$

- axion $a$ much heavier than $\phi$ during inflation, decouples:

  $m_\phi = \frac{M}{3}e^{-\sqrt{\frac{2}{3}}\phi_0} << m_a = \frac{M}{3}$

- inflation scale $M$ independent from NL-SUSY breaking scale $f$

  $\Rightarrow$ compatible with low energy SUSY

- string realization?

# Conclusions

String phenomenology:

Consistent framework for particle phenomenology and cosmology

possible 3 very different scales (besides $M_{Planck}$)

electroweak, dark energy, inflation

Maximal predictive power if common frame for:

moduli stabilization, model building, SUSY breaking and calculability

e.g. magnetized branes

- SUSY breaking with infinitesimal (tunable) +ve cosmological constant

  interesting framework for model building incorporating dark energy

- Inflation models at a hierarchically different third scale

  Sgoldstino-less supergravity models of inflation