28th International workshop "What comes beyond the standard models"

Stable multicharged particles in the early Universe

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Dark atom model

- Stable lepton-like multicharged heavy particles X^{-2n} arise in several models [1,2].
- Only even values of charge are allowed [3].
- Such particles should form neutral bound states with ordinary matter nuclei.

$$X^{-2n} + n He^{+2} \rightarrow XHe + \gamma$$

Outline

- Bound states and their description
 - Inner structure of dark ions
 - Estimations and applications
- Nuclear reactions with dark ions
 - Using analogies
 - Accurate calculation
- Primordial concentrations
 - Anomalous isotopes
 - Primordial metals
- Conclusions

Isotopes produced in primordial nucleosynthesis [4]:

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• What structure do dark XN ions have?

Inner structure of dark ions

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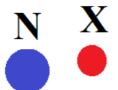


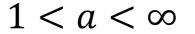
- Point-like particles
- ➤ Analogy with ordinary hydrogen:

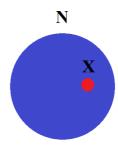
$$E_{X-N}^{\text{Bohr}} = 2n^2 Z^2 \alpha^2 m_N = \frac{1}{2m_N r_B^2}$$

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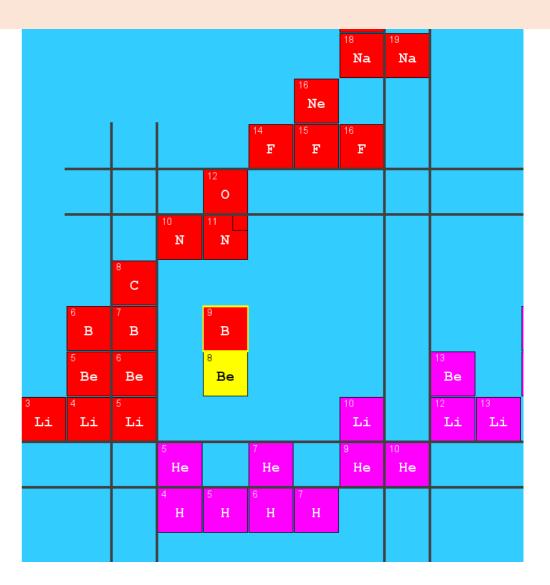
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- Non-point nucleus in a shell
- ➤ Nuclear analogies should be used.
- \triangleright Binding energies: $E_{X-N}^{Thomson}$ and Q_N

Inner structure of dark ions

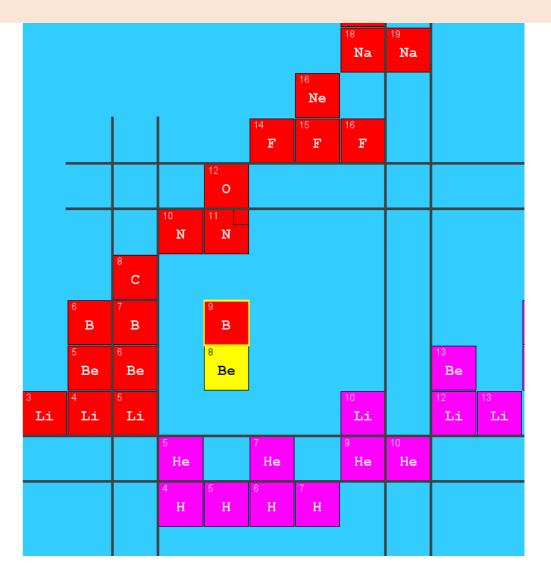
Isotope stabilization



Inner structure of dark ions

Isotope stabilization

• New stable-in-shall isotopes? $(Xp_4 \text{ for instance})$



The most states are Thomson-like.

• There are a lot of $E_{X-N}^{Thomson}$ estimations [5,6,7], but no Q_N estimations.

$$\mathcal{H} = egin{dcases} rac{p^2}{2m_N} - rac{Z_X Z_N lpha}{2r_N} \left(3 - rac{r^2}{r_N^2}
ight), & r < r_N \ rac{p^2}{2m_N} - rac{Z_X Z_N lpha}{r}, & r > r_N \end{cases}$$

$$E_{X-N}^{\text{Glashow}} = \begin{cases} \frac{3}{2} \frac{Z_X Z_N \alpha}{r_N} \left(1 - \sqrt{\frac{1}{a}} \right), & a \gtrsim 2\\ \frac{1}{m_N r_N} \lambda \left(a \right), & a \approx 1\\ \frac{1}{2} (Z_X Z_N \alpha)^2 m_N, & a < 1 \end{cases}$$

n	A			
	H	He	Li	Be
1	В	4	3	T
2	4	3	Т	Т
3	3	Т	Т	Т
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^[6] Akhmedov E. Nuclear fusion catalyzed by doubly charged scalars: Implications for energy production // Phys. Rev. D – 2022. – Vol. 106, No. 03

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- Reaction: $XN_1 \rightarrow XN_2 + N_3$
- Energy: $(Z_1m_p + (A_1-Z_1)m_n \Delta m_1^{shell}) (Z_2m_p + (A_2-Z_2)m_n \Delta m_2^{shell}) (Z_3m_p + (A_3-Z_3)m_n)$

Estimations and applications

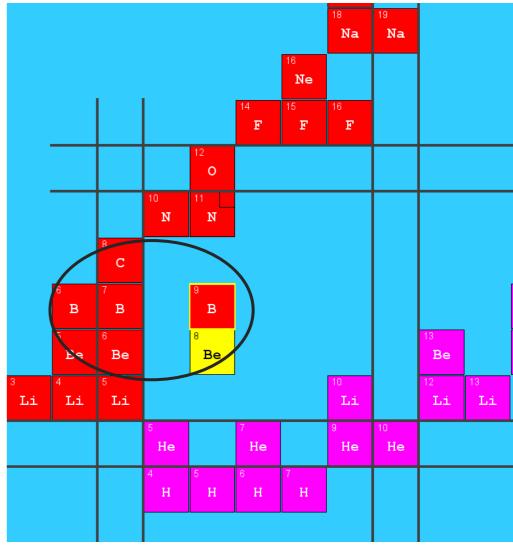
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• Prediction:

•
$$\begin{bmatrix} ^6Be, \ ^8Be \\ ^6B, \ ^7B, \ ^9B \end{bmatrix}$$
 are stabilizing $^8\mathcal{C}$

n-emission is not prevented



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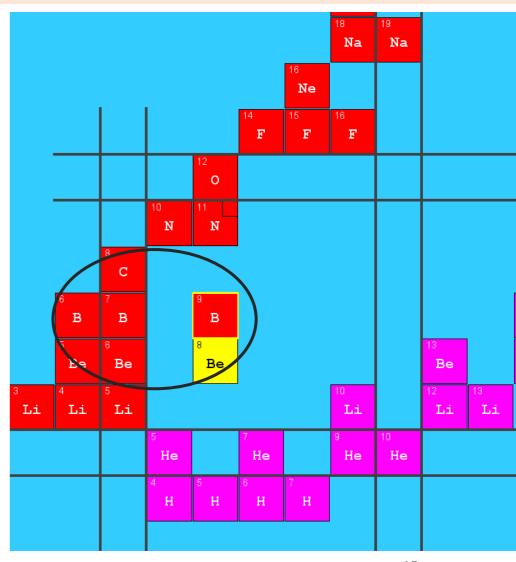
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- An error should be rather large

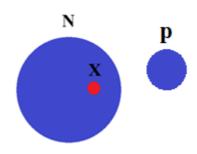


- Cluster approach:
 - Consider nucleus as a bound state of n, p, α clusters
 - Use **experimental** $V_{\alpha-\alpha}$, $V_{\alpha-N_{,}}$ and $V_{N-N_{,}}$ potentials to improve predictions for Li, Be, C isotopes

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 - \triangleright the energy of p-, α separation
 - > wave functions of bound states

How to describe XN more accurate?

Cluster approach

> Expected results:

- \triangleright the energy of p-, α separation
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➤ Potential problems:

- ➤ Is not applicable on high charges (n>1) because of non-point helium
- ➤ Bad predictions for spherical nuclei (???)

- Ab-initio approach [10,11]
 - Simulate nucleus by solving a many-body quantum problem
 - Use mesonic potential, modified by Coulomb term arising from the interaction with X^{-2n}

Nuclear reactions with dark ions

$$(1) X + N \to XN + \gamma$$

(2)
$$XN_1 + N_2 \rightarrow XN_3 + N_4/\gamma$$

(3)
$$XN_1 + XN_2 \rightarrow X_2N_3 + N_4/\gamma$$

- (1) The first stage of dark atom recombination
- (2) Nuclear interaction of ordinary matter with dark ions
- (3) Dark chemistry: interaction of two dark ions

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What does we need?

- ➤ List (network) of reactions
- ➤ Cross sections/rates

(1)
$$X + N \rightarrow XN + \gamma$$

Bohr states (point-like nucleus!):
 Recombination of hydrogen-like atom [13]

The exact analogy can be used for capturing of:

$$\circ$$
 n=1: $p, d, t, {}^{3}He$

$$\langle \sigma v \rangle_{\text{rec}} = \frac{32}{3} \sqrt{\frac{2\pi E_{X-N}}{3 m_N}} (Z_X Z_N \alpha)^3 r_B^2 \cdot \left(\frac{E_{X-N}}{T}\right)^{\frac{1}{2}} \left(\ln \left(\frac{E_{X-N}}{T}\right) + \gamma\right)$$

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Thomson states (non-point-like nucleus)
 The picture changes completely.
 it is impossible to use the hydrogen-like wave function, and therefore the cross section must change.

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(2)
$$XN_1 + N_2 \rightarrow XN_3 + N_4/\gamma$$

• If the similar ordinary reaction is exist [4]

$$\langle \sigma v \rangle^{ij} (T) = \frac{2}{T} \sqrt{\frac{2}{\pi T} \frac{m_i + m_j}{m_i m_j}} \int_0^\infty \sigma^{ij}(E) E e^{-\frac{E}{T}} dE$$

$$\sigma^{ij}(E) = \frac{S^{ij}(E)}{E} e^{-2\pi\eta_{ij}},$$

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In [7] the accurate calculations for X^- [14] was rescaled. Is it any difference and why?

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- If not:
 - rough analogies [15]:

$$(XN_1)^0 + N_2^{+q} \to (XN_3)^{+q} + \gamma$$
 and $n + p \to d + \gamma$
$$\langle \sigma v \rangle = \frac{f\pi\alpha}{m_p^2} \frac{3}{\sqrt{2}} \left(\frac{Z}{A}\right)^2 \frac{T}{\sqrt{Am_p E}},$$

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$$OHe + He \rightarrow OBe + \gamma$$

Nuclear reactions with dark ions Accurate calculation

- Accurate calculation of reaction rates [6,14] requires to know wave functions of particles
- To find it
 - In [6] the variation method was used
 - In [14] the quantum three-body problem was solved

1)
$$O^{--} + {}^{4}\text{He} \rightarrow OHe + \gamma$$
;

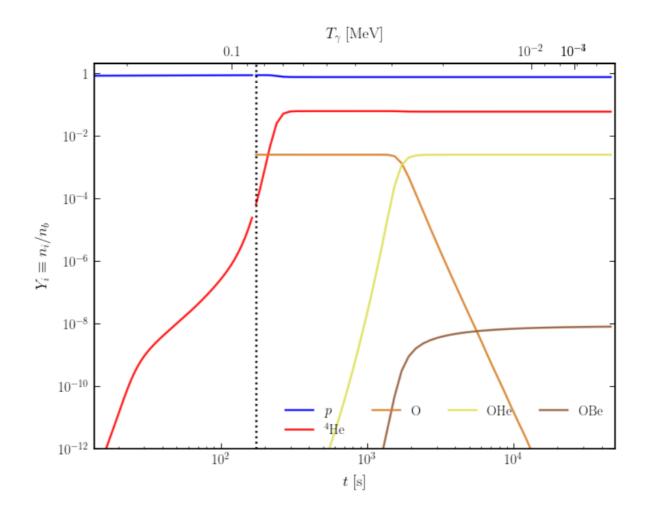
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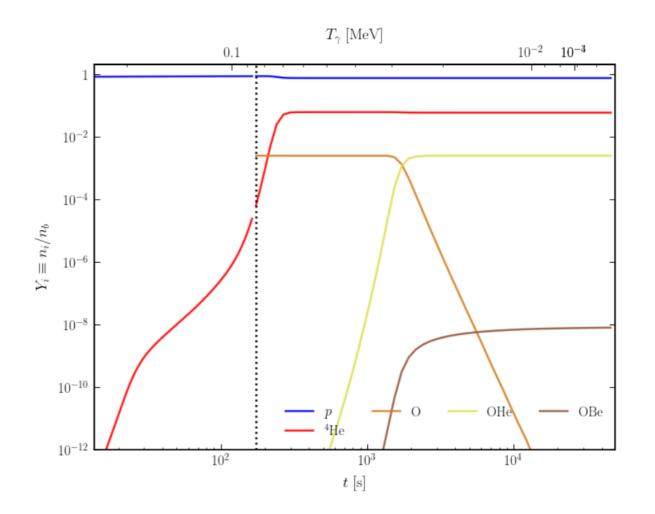


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$$\begin{split} \left\langle \sigma v \right\rangle_{\rm rec} &= \frac{32}{3} \sqrt{\frac{2\pi E_{X-N}}{3 \, m_N}} \left(Z_X Z_N \, \alpha \right)^3 r_B^2 \cdot \left(\frac{E_{X_N}}{T} \right)^{\frac{1}{2}} \left(\ln \left(\frac{E_{X_N}}{T} \right) + \gamma \right), \\ \left\langle \sigma v \right\rangle &= \frac{f \pi \alpha}{m_p^2} \frac{3}{\sqrt{2}} \left(\frac{Z}{A} \right)^2 \frac{T}{\sqrt{A m_p E}}, \end{split}$$

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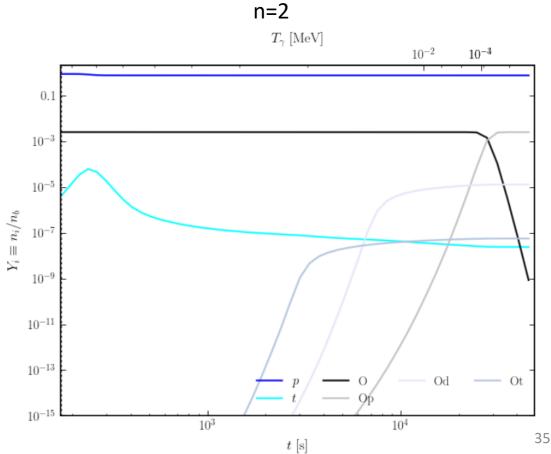


Exact rates are known only for capturing of:

o n=1: $p, d, t, {}^{3}He$

o n=2: *p, d, t*

n=1 $T_{\gamma} [\text{MeV}]$ 10^{-2} 10^{-3} 0.1 10^{-3} 10^{-5} $Y_i \equiv n_i/n_b$ 10^{-11} 10^{-13} 10^{-13} 10^{3} 10^{4}



$$T_{\text{rec}} = E_{X-N} \left(\ln \left(\frac{g_X g_N}{g_{XN}} \left(\frac{m_N T_{\text{now}}^2}{2\pi E_{X-N}} \right)^{\frac{3}{2}} \frac{1}{n_N^{\text{now}}} \right) \right)^{-1}$$

$$E_{X-N}^{\text{Bohr}} = 2n^2 Z^2 \alpha^2 m_N = \frac{1}{2m_N r_D^2}$$

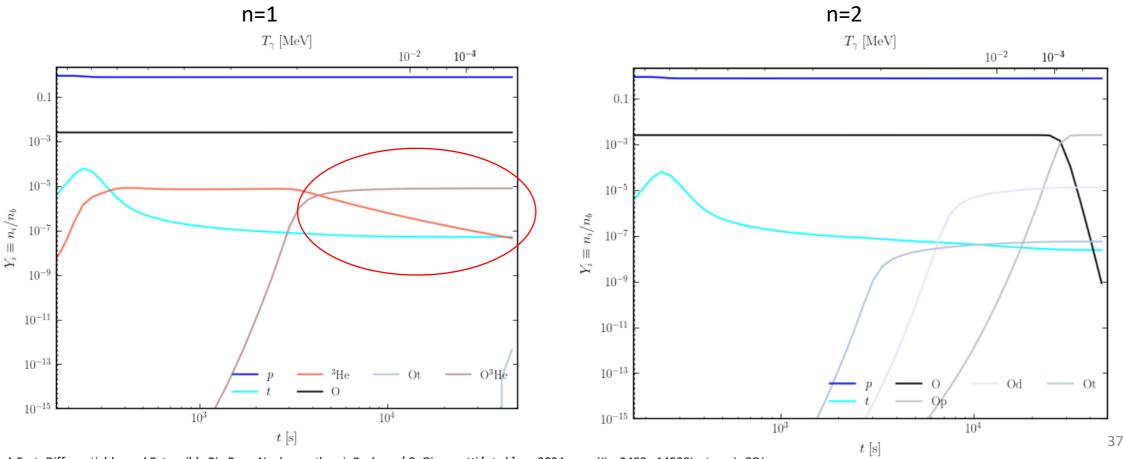
$$E_{W-N}^{\text{Thompson}} = \frac{3}{2} \frac{Z_X Z_N \alpha}{r_N} \left(1 - \sqrt{\frac{r_B}{r_N}} \right)$$

	$T_{ m rec},$ кэ ${ m B}$				
$ ^n $	p	D	3He	4He	
1	3 (B)	4 (B)	28 (B)	37 (B)	
2	13 (B)	19 (B)	~ 42 (T)	85 (T)	
3	29 (B)	44 (B)	$\sim 116 \; (T)$	180 (T)	
4	54 (B)	79 (B)	$\sim 198 \; (T)$	285 (T)	
5	86 (B)	$\sim 126 \; (B)$	$\sim 286 \; (T)$	395 (T)	

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[16] LINX: A Fast, Differentiable, and Extensible Big Bang Nucleosynthe-sis Package / C. Giovanetti [et al.]. — 2024. — arXiv: 2408 . 14538[astro-ph.CO].

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- Gamma-ray spectrum analysis?

Conclusions

- To predict the primordial concentrations of dark ions, it is necessary to describe both the reactions of nuclear shell formation and the internal structure:
 - kinetic equations require reaction rates, which for the most part can be found by rescaling
 - Both binding energies E_{X-N} and Q_N are needed to create an accurate network of reactions
 - The nuclear shell can be described using the methods of modern nuclear physics.
- The overproduction of anomalous isotopes and primordial metals may constrain the dark atom model, but the standard model of nucleosynthesis should be revised

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